

The Topology of Circuit-Field Coupling

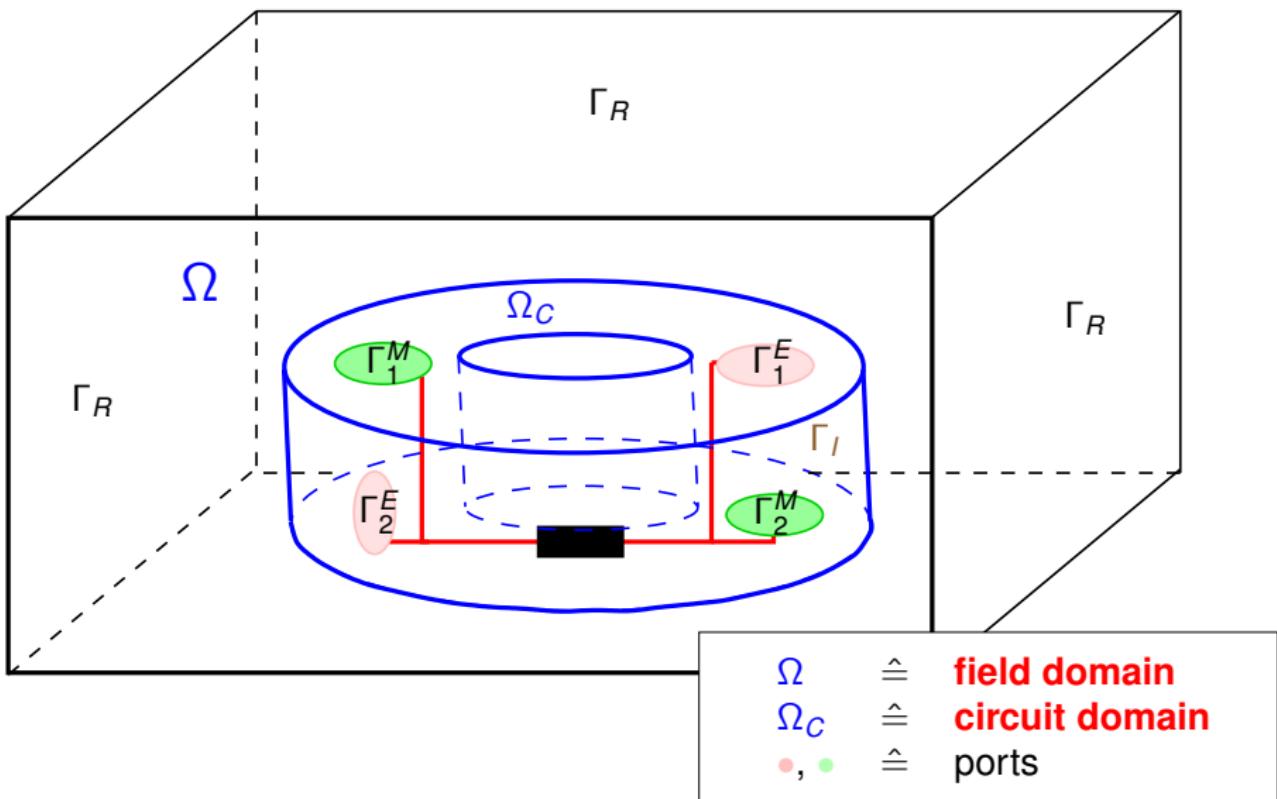
R. Hiptmair¹, J. Ostrowski²

¹ Seminar for Applied Mathematics, ETH Zürich

² ABB Switzerland Ltd., Corporate Research, Segelhofstrasse 1, CH-5405 Baden

Mathematics Colloquium
University of Illinois at Urbana-Champaign

Circuit Domain Ω_C – Field Domain Ω



Coupling interface:

$$\Gamma := \partial\Omega_C = \Gamma_1^E \cup \dots \cup \Gamma_{N_E}^E \cup \Gamma_1^M \cup \dots \cup \Gamma_{N_M}^M \cup \Gamma_I$$

Maxwell's Equations

$\Sigma \triangleq$ oriented surface

Faraday's law

$$\int_{\partial\Sigma} \mathbf{E} \cdot d\mathbf{s} = -\frac{d}{dt} \int_{\Sigma} \mathbf{B} \cdot \mathbf{n} dS$$



$$\text{curl } \mathbf{E} = -\partial_t \mathbf{B}$$

Ampere's law

$$\int_{\partial\Sigma} \mathbf{H} \cdot d\mathbf{s} = \frac{d}{dt} \int_{\Sigma} \mathbf{D} \cdot \mathbf{n} dS$$



$$\text{curl } \mathbf{H} = \partial_t \mathbf{D}$$

Material laws:

$$\mathbf{B} = \mu(\mathbf{x})\mathbf{H} \quad , \quad \mathbf{D} = \epsilon(\mathbf{x})\mathbf{E}$$

$\left. \begin{matrix} \mathbf{E}, \mathbf{H} \\ \mathbf{D}, \mathbf{B} \end{matrix} \right\}$ integrated over $\left\{ \begin{matrix} \text{paths} \\ \text{surfaces} \end{matrix} \right\}$



Exterior calculus perspective : $\begin{matrix} \mathbf{E}, \mathbf{H} & \leftrightarrow & 1\text{-forms} \\ \mathbf{D}, \mathbf{B} & \leftrightarrow & 2\text{-forms} \\ \text{curl} & \leftrightarrow & \text{exterior derivative } d \end{matrix}$

Maxwell's equations



De Rham complex

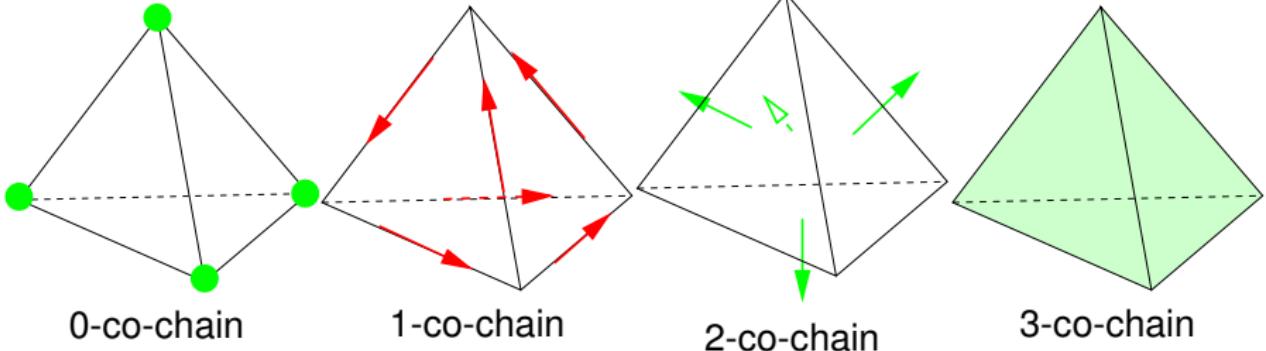


algebraic topology

Discrete Fields

$\Omega_h \hat{=} \text{oriented tetrahedral mesh for } \Omega \text{ (simplicial complex)}$

(Discrete) ℓ -co-chain $\vec{U} \in \mathcal{C}^\ell(\Omega_h)$: $\vec{U} : \{ \ell\text{-facets of } \Omega_h \} \mapsto \mathbb{R}$



$$\begin{aligned}\mathbf{E}(t), \mathbf{H}(t) &\leftrightarrow 1\text{-co-chains} \quad \vec{\mathbf{E}}(t), \vec{\mathbf{H}}(t) \in \mathcal{C}^1(\Omega_h) \quad (\text{d.o.f.s on edges}) \\ \mathbf{B}(t), \mathbf{D}(t) &\leftrightarrow 2\text{-co-chains} \quad \vec{\mathbf{B}}(t), \vec{\mathbf{D}}(t) \in \mathcal{C}^2(\Omega_h) \quad (\text{d.o.f.s on faces})\end{aligned}$$

► $\left\{ \begin{array}{l} \text{Faraday's law: } \mathbf{C}\vec{\mathbf{E}} = -\partial_t \vec{\mathbf{B}} \\ \text{Ampere's law: } \mathbf{C}\vec{\mathbf{H}} = \partial_t \vec{\mathbf{D}} \end{array} \right. , \quad \mathbf{C} \hat{=} \begin{array}{l} \text{edge-face} \\ \text{incidence matrix} \end{array}$

Polynomial extension of ℓ -cochains ► Whitney- ℓ -forms ► FEEC

Maxwell's Equations: Weak Formulations

I. **E-based**: I.b.p. on Ampere's law & keep Faraday's law

$$\begin{aligned} \int_{\Omega} \partial_t(\epsilon \mathbf{E}) \cdot \mathbf{E}' - \mathbf{H} \cdot \operatorname{curl} \mathbf{E}' \, d\mathbf{x} - \int_{\Gamma} (\mathbf{H} \times \mathbf{n}) \cdot \mathbf{E}' \, dS &= 0 \quad \forall \mathbf{E}', \\ \int_{\Omega} \partial_t(\mu \mathbf{H}) \cdot \mathbf{H}' + \operatorname{curl} \mathbf{E} \cdot \mathbf{H}' \, d\mathbf{x} &= 0 \quad \forall \mathbf{H}'. \end{aligned}$$

Spaces: $\mathbf{E}(t) \in \mathbf{H}(\operatorname{curl}, \Omega) := \{\mathbf{V} \in \mathbf{L}^2(\Omega) : \operatorname{curl} \mathbf{V} \in \mathbf{L}^2(\Omega)\}$, $\mathbf{H}(t) \in \mathbf{L}^2(\Omega)$

II. **H-based**: I.b.p. on Faraday's law & keep Ampere's law

$$\begin{aligned} \int_{\Omega} \partial_t(\mu \mathbf{H}) \cdot \mathbf{H}' - \mathbf{E} \cdot \operatorname{curl} \mathbf{H}' \, d\mathbf{x} - \int_{\Gamma} \mathbf{E} \cdot (\mathbf{H}' \times \mathbf{n}) \, dS &= 0 \quad \forall \mathbf{H}', \\ \int_{\Omega} \partial_t(\epsilon \mathbf{E}) \cdot \mathbf{E}' + \operatorname{curl} \mathbf{H} \cdot \mathbf{E}' \, d\mathbf{x} &= 0 \quad \forall \mathbf{E}'. \end{aligned}$$

Spaces: $\mathbf{H}(t) \in \mathbf{H}(\operatorname{curl}, \Omega) := \{\mathbf{V} \in \mathbf{L}^2(\Omega) : \operatorname{curl} \mathbf{V} \in \mathbf{L}^2(\Omega)\}$, $\mathbf{E}(t) \in \mathbf{L}^2(\Omega)$

Traces

$$\int_{\Omega} \partial_t(\epsilon \mathbf{E}) \cdot \mathbf{E}' - \mathbf{H} \cdot \operatorname{curl} \mathbf{E}' \, d\mathbf{x} - \int_{\Gamma} \gamma_x \mathbf{H} \cdot \gamma_t \mathbf{E}' \, dS = 0 \quad \forall \mathbf{E}',$$

$$\int_{\Omega} \partial_t(\mu \mathbf{H}) \cdot \mathbf{H}' + \operatorname{curl} \mathbf{E} \cdot \mathbf{H}' \, d\mathbf{x} = 0 \quad \forall \mathbf{H}'.$$

γ_p	$\hat{=}$	pointwise trace	$\gamma_p U(\mathbf{x}) := U(\mathbf{x}),$
γ_t	$\hat{=}$	tangential trace	$\gamma_t \mathbf{U}(\mathbf{x}) := \mathbf{n}(\mathbf{x}) \times (\mathbf{U}(\mathbf{x}) \times \mathbf{n}(\mathbf{x})), \quad \mathbf{x} \in \Gamma.$
γ_x	$\hat{=}$	twisted tangential trace	$\gamma_x \mathbf{U}(\mathbf{x}) := \mathbf{U}(\mathbf{x}) \times \mathbf{n}(\mathbf{x}),$
γ_n	$\hat{=}$	normal trace	$\gamma_n \mathbf{U}(\mathbf{x}) := \mathbf{U}(\mathbf{x}) \cdot \mathbf{n}(\mathbf{x}),$

$$\begin{array}{ccccccc}
 H^1(\Omega) & \xrightarrow{\text{grad}} & \boldsymbol{H}(\operatorname{curl}, \Omega) & \xrightarrow{\text{curl}} & \boldsymbol{H}(\operatorname{div}, \Omega) & \xrightarrow{\text{div}} & L^2(\Omega) \\
 \gamma_p \downarrow & & \gamma_t \downarrow & & \gamma_n \downarrow & & \\
 H^{\frac{1}{2}}(\Gamma) & \xrightarrow{\text{grad}_\Gamma} & \boldsymbol{H}^{-\frac{1}{2}}(\operatorname{curl}_\Gamma, \Gamma) & \xrightarrow{\text{curl}_\Gamma} & H^{-\frac{1}{2}}(\Gamma) & &
 \end{array}$$

Traces

γ_p	$\hat{=}$	pointwise trace	$\gamma_p U(\mathbf{x}) := U(\mathbf{x}),$
γ_t	$\hat{=}$	tangential trace	$\gamma_t \mathbf{U}(\mathbf{x}) := \mathbf{n}(\mathbf{x}) \times (\mathbf{U}(\mathbf{x}) \times \mathbf{n}(\mathbf{x})), \quad \mathbf{x} \in \Gamma.$
γ_{\times}	$\hat{=}$	twisted tangential trace	$\gamma_{\times} \mathbf{U}(\mathbf{x}) := \mathbf{U}(\mathbf{x}) \times \mathbf{n}(\mathbf{x}),$
γ_n	$\hat{=}$	normal trace	$\gamma_n \mathbf{U}(\mathbf{x}) := \mathbf{U}(\mathbf{x}) \cdot \mathbf{n}(\mathbf{x}),$

$$\begin{array}{ccccccc}
 H^1(\Omega) & \xrightarrow{\text{grad}} & \boldsymbol{H}(\mathbf{curl}, \Omega) & \xrightarrow{\text{curl}} & \boldsymbol{H}(\text{div}, \Omega) & \xrightarrow{\text{div}} & L^2(\Omega) \\
 \gamma_p \downarrow & & \gamma_t \downarrow & & \gamma_n \downarrow & & \\
 H^{\frac{1}{2}}(\Gamma) & \xrightarrow{\text{grad}_{\Gamma}} & \boldsymbol{H}^{-\frac{1}{2}}(\mathbf{curl}_{\Gamma}, \Gamma) & \xrightarrow{\text{curl}_{\Gamma}} & H^{-\frac{1}{2}}(\Gamma) & &
 \end{array}$$

$\Gamma_h := \Omega_h|_{\Gamma}$ = mesh of interface Γ !

Trace of ℓ -co-chain $\vec{V} \in \mathcal{C}^\ell(\Omega_h)$: $\gamma \vec{V} := \vec{V} \Big|_{\Gamma_h}$

$$\begin{array}{ccccccc}
 \mathcal{C}^0(\Omega_h) & \xrightarrow{\text{G}} & \mathcal{C}^1(\Omega_h) & \xrightarrow{\text{C}} & \mathcal{C}^2(\Omega_h) & \xrightarrow{\text{D}} & \mathcal{C}^3(\Omega_h) \\
 \gamma \downarrow & & \gamma \downarrow & & \gamma \downarrow & & \\
 \mathcal{C}^0(\Gamma_h) & \xrightarrow{\text{G}_{\Gamma}} & \mathcal{C}^1(\Gamma_h) & \xrightarrow{\text{C}_{\Gamma}} & \mathcal{C}^2(\Gamma_h) & &
 \end{array}$$

Coupling Conditions and Ports

Partition: $\Gamma := \partial\Omega_C = \underbrace{\Gamma_1^E \cup \dots \cup \Gamma_{N_E}^E}_{=: \Gamma_E} \cup \underbrace{\Gamma_1^M \cup \dots \cup \Gamma_{N_M}^M}_{=: \Gamma_M} \cup \Gamma_I$

electric ports magnetic ports insulating boundary

@*electric ports*: $\gamma_t \mathbf{E} = 0$ on Γ_E (PEC)

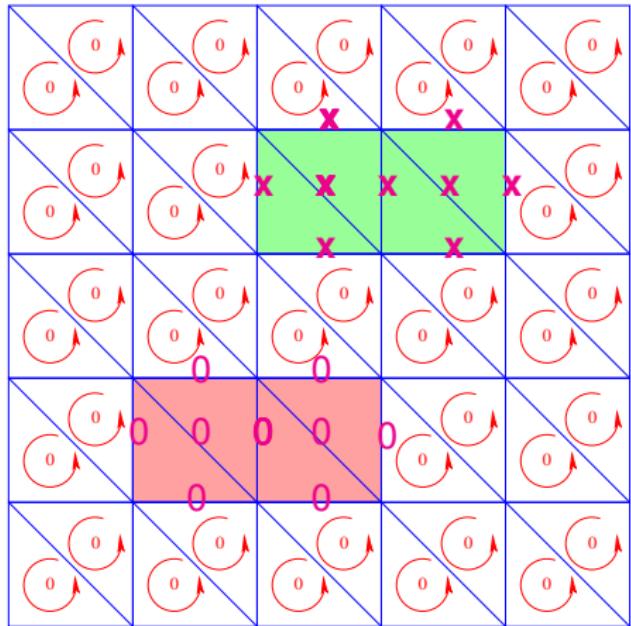
free electric current no magnetic flux

@*magnetic ports*: $\gamma_x \mathbf{H} = 0$ on Γ_M (PMC)

no electric current free magnetic flux

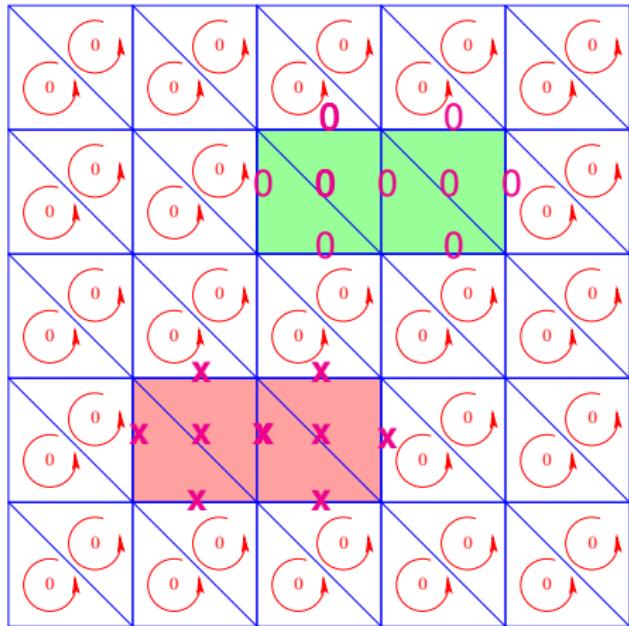
@*insulating boundary*: no magnetic flux: $\text{curl } \mathbf{E} \cdot \mathbf{n} = \text{curl}_\Gamma(\gamma_t \mathbf{E}) = 0$,
no electric current: $\text{curl } \mathbf{H} \cdot \mathbf{n} = \text{div}_\Gamma(\gamma_x \mathbf{H}) = 0$.

Discrete Surface Fields



$$\gamma \vec{\mathbf{E}} \in \mathcal{C}^1(\Gamma_h)$$

$\hat{=} \Gamma_E$, $\hat{=} \Gamma_M$, 0 $\hat{=}$ zero edge circulation, X $\hat{=}$ free edge value

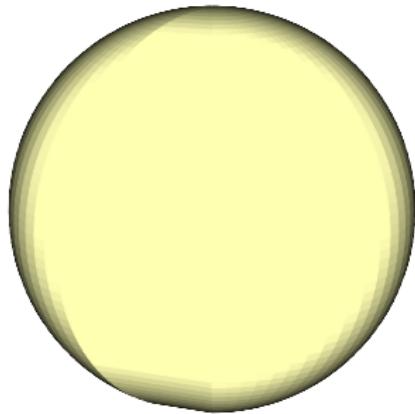


$$\gamma \vec{\mathbf{H}} \in \mathcal{C}^1(\Gamma_h)$$

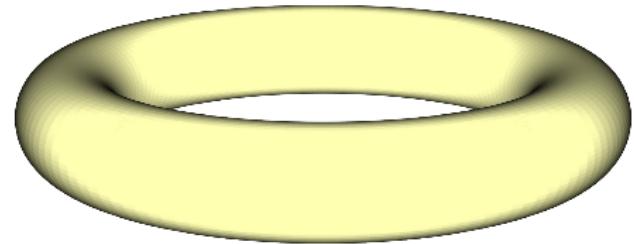
Surface Scalar Potentials

If Ω has no handles $\Leftrightarrow \beta_1(\Omega) = \beta_1(\Gamma) = 0$, then

$$\mathbf{v} \in \mathbf{H}^{-\frac{1}{2}}(\operatorname{curl}_\Gamma, \Gamma), \quad \operatorname{curl}_\Gamma \mathbf{v} = 0 \quad \Rightarrow \quad \exists \varphi \in H^{\frac{1}{2}}(\Gamma): \quad \mathbf{v} = \operatorname{grad}_\Gamma \varphi,$$
$$\vec{v} \in \mathcal{C}^1(\Gamma_h), \quad \mathbf{C}_\Gamma \vec{v} = 0 \quad \Rightarrow \quad \exists \vec{\varphi} \in \mathcal{C}^0(\Gamma_h): \quad \vec{v} = \mathbf{G}_\Gamma \vec{\varphi}.$$



$$\beta_1(\Gamma) = 0$$



$$\beta_1(\Gamma) = 1$$

Surface Electric Fields

Setting: Electric ports Γ_E , magnetic ports Γ_M & insulating boundary Γ_I

$$\gamma_t \mathbf{E} \in \mathcal{V}_E(\Gamma) := \left\{ \mathbf{v} \in \mathbf{H}^{-\frac{1}{2}}(\operatorname{curl}_{\Gamma}, \Gamma) : \operatorname{curl}_{\Gamma} \mathbf{v} = 0 \text{ in } \Gamma \setminus \Gamma_M, \mathbf{v}|_{\Gamma_E} = 0 \right\}.$$

► $\forall \mathbf{u} \in \mathcal{V}_E(\Gamma): \int_{\sigma} \mathbf{u} \cdot d\vec{s} = 0 \quad \forall \sigma \in \mathcal{B}(\Gamma \setminus \Gamma_M),$

$$\mathcal{B}(\Gamma \setminus \Gamma_M) := \{\text{boundaries in } \Gamma \setminus \Gamma_M\}.$$

Space for scalar potentials:

$$\mathcal{S}(\Gamma) := \left\{ \varphi \in H^{\frac{1}{2}}(\Gamma \setminus \Gamma_M) : \varphi = 0 \text{ on } \Gamma_E \right\}.$$

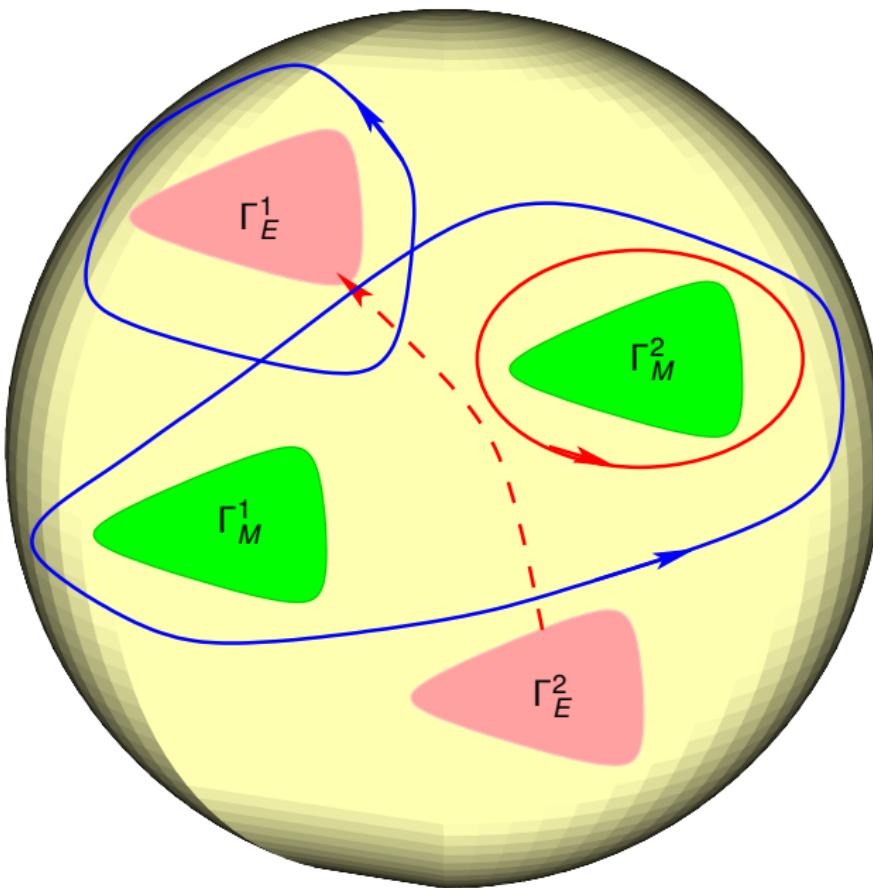
► $\forall \mathbf{u} \in \operatorname{grad}_{\Gamma} \mathcal{S}(\Gamma): \int_{\sigma} \mathbf{u} \cdot d\vec{s} = 0 \quad \forall \sigma \in \mathcal{Z}(\Gamma \setminus \Gamma_M; \partial \Gamma_E),$

$$\mathcal{Z}(\Gamma \setminus \Gamma_M; \Gamma_E) := \{\text{paths either closed or with endpoints } \in \Gamma_E\}.$$

$$\mathcal{B}(\Gamma \setminus \Gamma_M) \subset \mathcal{Z}(\Gamma \setminus \Gamma_M; \partial \Gamma_E),$$

but $\mathcal{V}_E(\Gamma) \neq \operatorname{grad}_{\Gamma} \mathcal{S}(\Gamma)$, if $\mathcal{B}(\Gamma \setminus \Gamma_M) \neq \mathcal{Z}(\Gamma \setminus \Gamma_M; \Gamma_E)$.

Boundaries and Cycles



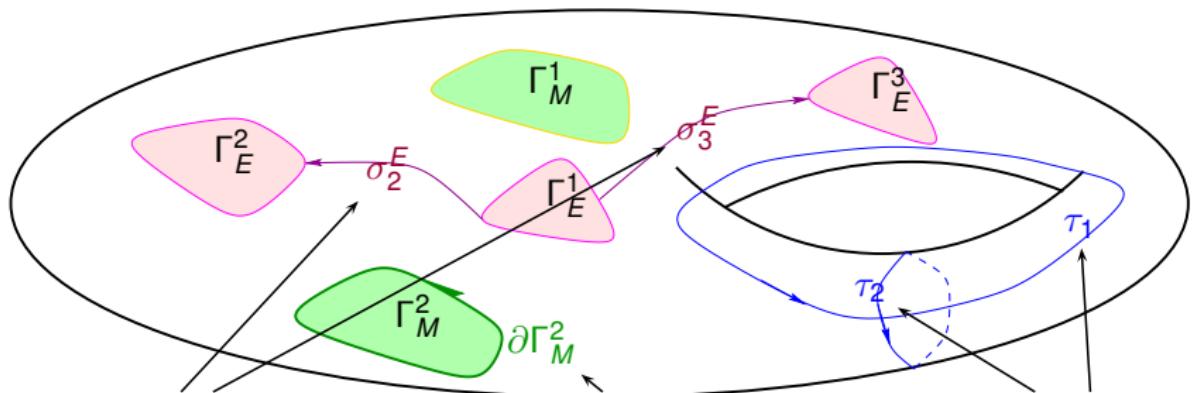
$\mathcal{B}(\Gamma \setminus \Gamma_M) := \{$
boundaries in $\Gamma \setminus \Gamma_M\},$

$\mathcal{Z}(\Gamma \setminus \Gamma_M; \Gamma_E) := \{$
paths either closed or
with endpoints $\in \Gamma_E\} .$

- $\in \mathcal{Z}(\Gamma \setminus \Gamma_M; \Gamma_E)$
 $\notin \mathcal{B}(\Gamma \setminus \Gamma_M)$
- $\in \mathcal{B}(\Gamma \setminus \Gamma_M)$
- - - $\in \mathcal{Z}(\Gamma \setminus \Gamma_M; \Gamma_E)$

Fundamental Relative Cycles

$$\mathcal{Z}(\Gamma \setminus \Gamma_M; \Gamma_E) = \mathcal{B}(\Gamma \setminus \Gamma_M) + \text{span}\{M \text{ fundamental cycles}\} \quad .$$



► electric connector paths

magnetic port cycles

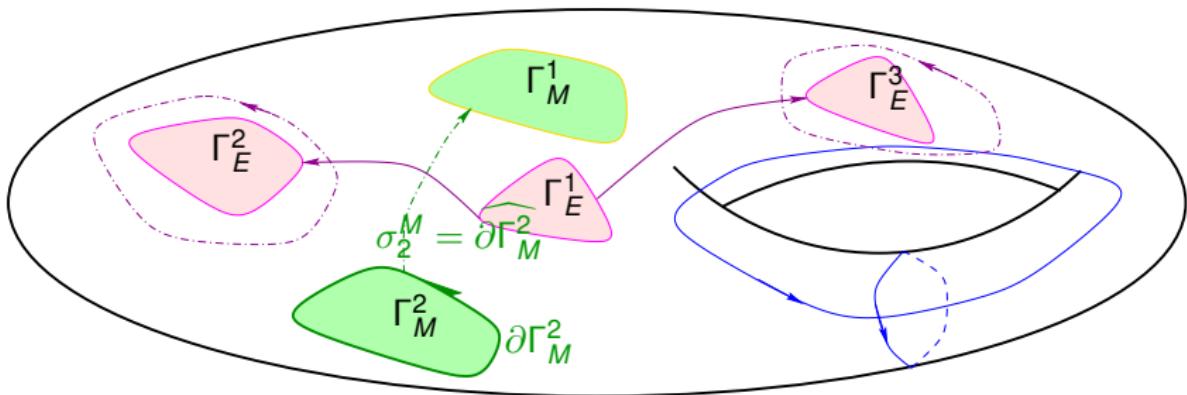
topological cycles

$$M = \max\{N_E, 1\} - 1 + \max\{N_M, 1\} - 1 + 2\beta_1(\Omega) \quad .$$

$\{\phi_1, \dots, \phi_M\} \hat{=} \text{fundamental cycles of } \mathcal{Z}/\mathcal{B}$. $\exists \mathbf{c}_1, \dots, \mathbf{c}_M \in \mathcal{V}_E(\Gamma)$:

$$\int_{\phi_j} \mathbf{c}_i \cdot d\mathbf{s} = \delta_{ij} \quad , \quad \mathcal{V}_E(\Gamma) = \mathbf{grad}_{\Gamma} S(\Gamma) + \text{span}\{\mathbf{c}_1, \dots, \mathbf{c}_M\} \quad .$$

Tool: Dual Relative Cycles



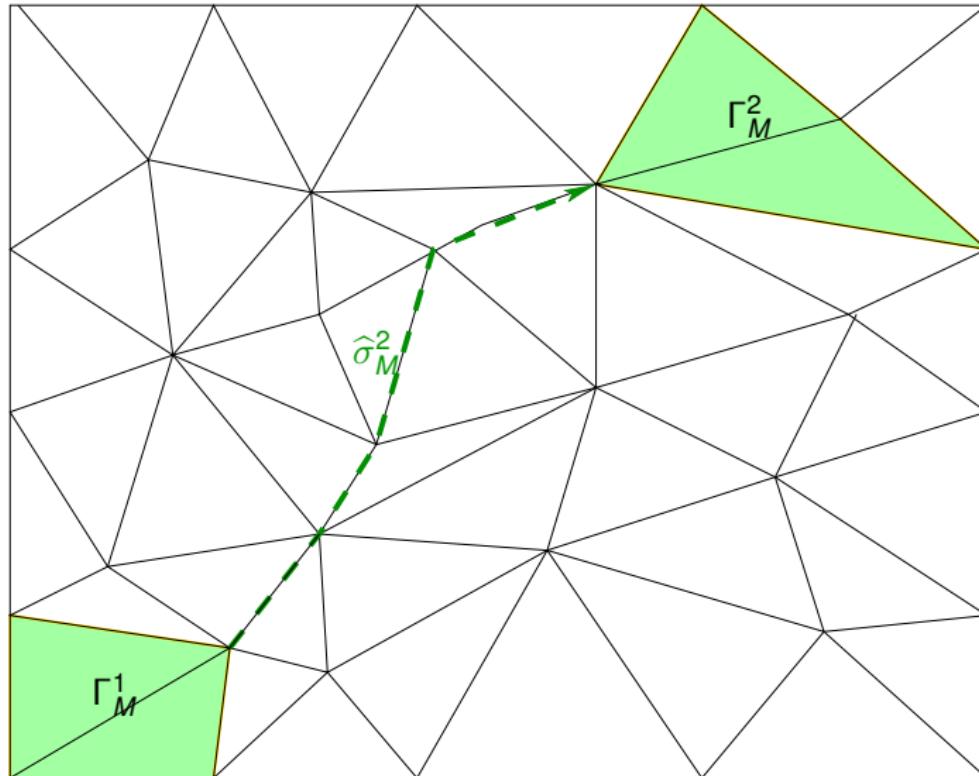
Poincaré
Lefschetz
duality

	fundamental cycle	dual cycle
	electric connector σ_k^E	electric port cycle $\hat{\sigma}_k^E$
	magnetic port cycle σ_ℓ^M	magnetic connector $\hat{\sigma}_\ell^M$

|

Collar Fields

fundamental cycle σ \longrightarrow dual cycle $\widehat{\sigma}$ \longrightarrow collar field $\mathbf{c} \in \mathcal{V}_E(\Gamma)$



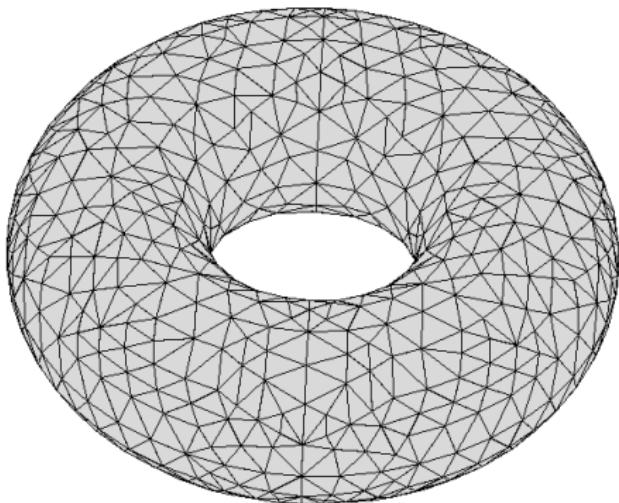
$\vec{\mathbf{c}} \in \mathcal{C}_E^1(\Gamma_h)$:
(magnetic connector)

Fundamental Cycles: Construction



Spanning-tree techniques

Example: Torus ($\beta_1(\Gamma) = 2$) ► 2 topological fundamental cycles τ_1, τ_2

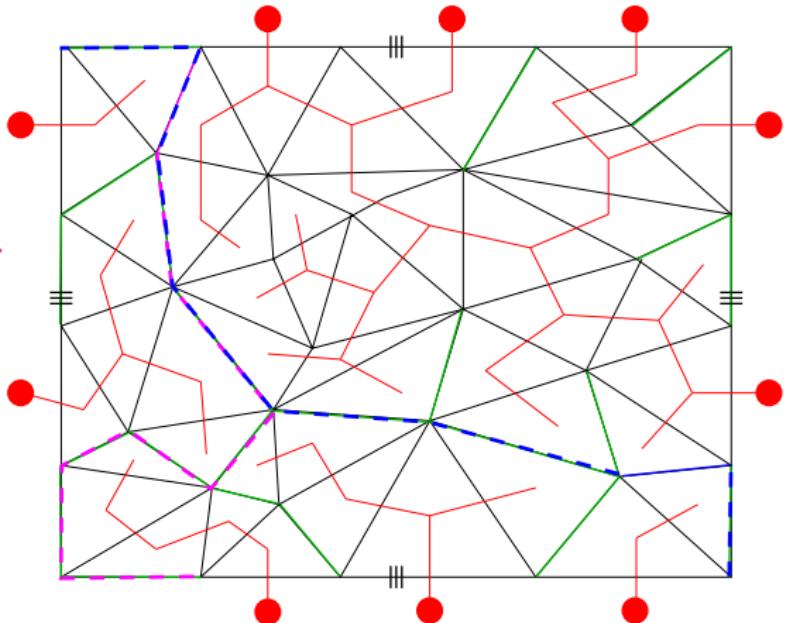


Fundamental Cycles: Construction

Demonstration:

Spanning-tree techniques

- : spanning tree in edge-cell graph
- : spanning tree in *remainder* edge-vertex graph
- , —: “Belt buckles”
- , —: fundamental cycles τ_1, τ_2

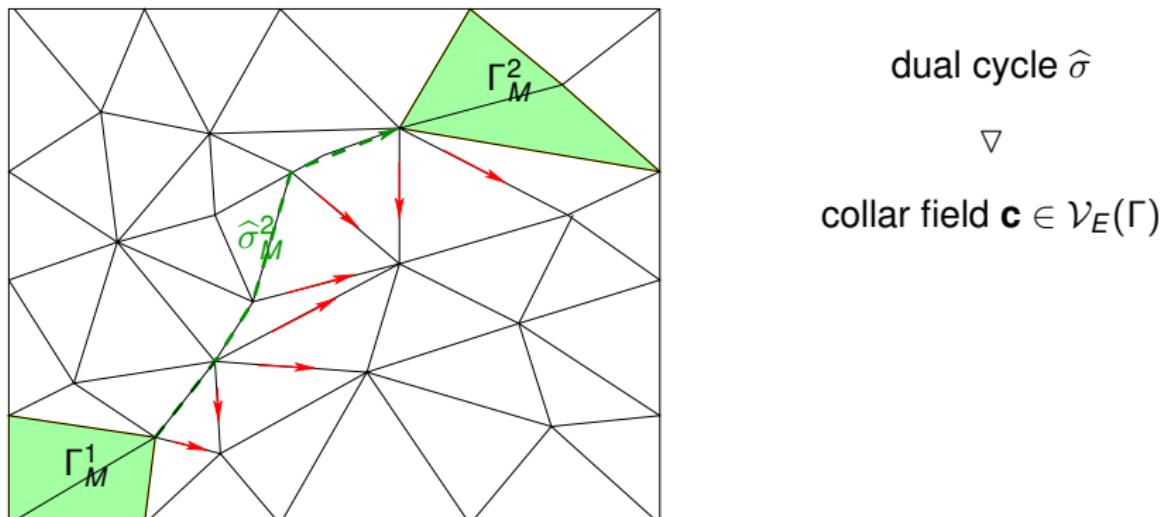


Topology → Variational Formulation

$$\int_{\Omega} \partial_t(\epsilon \mathbf{E}) \cdot \mathbf{E}' - \mathbf{H} \cdot \operatorname{curl} \mathbf{E}' \, d\mathbf{x} - \int_{\Gamma} \gamma_x \mathbf{H} \cdot \gamma_t \mathbf{E}' \, dS = 0 \quad \forall \mathbf{E}',$$

$$\int_{\Omega} \partial_t(\mu \mathbf{H}) \cdot \mathbf{H}' + \operatorname{curl} \mathbf{E} \cdot \mathbf{H}' \, d\mathbf{x} = 0 \quad \forall \mathbf{H}'.$$

$$\gamma_t \mathbf{E}' = \operatorname{grad} \varphi + \sum_{k=2}^{N_E} \mu_k \mathbf{c}_k^E + \sum_{\ell=2}^{N_M} \alpha_\ell \mathbf{c}_\ell^M + \sum_{m=1}^{N_T} \beta_m \mathbf{c}_m^T, \quad \varphi \in \mathcal{S}(\Gamma), \\ \mu_k, \alpha_\ell, \beta_m \in \mathbb{R}.$$



Topology \rightarrow Variational Formulation

$$\int_{\Omega} \partial_t(\epsilon \mathbf{E}) \cdot \mathbf{E}' - \mathbf{H} \cdot \operatorname{curl} \mathbf{E}' \, d\mathbf{x} - \int_{\Gamma} \gamma_x \mathbf{H} \cdot \gamma_t \mathbf{E}' \, dS = 0 \quad \forall \mathbf{E}',$$

$$\int_{\Omega} \partial_t(\mu \mathbf{H}) \cdot \mathbf{H}' + \operatorname{curl} \mathbf{E} \cdot \mathbf{H}' \, d\mathbf{x} = 0 \quad \forall \mathbf{H}'.$$

$$\gamma_t \mathbf{E}' = \mathbf{grad} \varphi + \sum_{k=2}^{N_E} \mu_k \mathbf{c}_k^E + \sum_{\ell=2}^{N_M} \alpha_\ell \mathbf{c}_\ell^M + \sum_{m=1}^{N_T} \beta_m \mathbf{c}_m^T, \quad \varphi \in \mathcal{S}(\Gamma),$$

$$\mu_k, \alpha_\ell, \beta_m \in \mathbb{R}.$$

Collar field $\mathbf{c} = \widetilde{\mathbf{grad}}_{\Gamma} \psi, \quad \psi \in H^1(\Gamma \setminus (\Gamma_M \cup \sigma)), \quad [\![\psi]\!]_{\widehat{\sigma}} = 1.$

►
$$\int_{\Gamma} \gamma_x \mathbf{H} \cdot \mathbf{c} \, dS = \int_{\Gamma \setminus \widehat{\sigma}} \gamma_x \mathbf{H} \cdot \mathbf{grad}_{\Gamma} \psi \, dS$$

$$= - \underbrace{\int_{\Gamma_I} \operatorname{div}_{\Gamma}(\gamma_x \mathbf{H}) \psi \, dS}_{=0} + \int_{\widehat{\sigma}} [\![\psi]\!]_{\widehat{\sigma}} \mathbf{H} \cdot d\mathbf{s} = \int_{\widehat{\sigma}} \mathbf{H} \cdot d\mathbf{s}.$$

Port Quantities

$$\mathbf{E} = \mathbf{E}_0 + \mathbf{grad} \Phi + \sum_{k=2}^{N_E} \mathbf{U}_k(t) \mathbf{C}_k^E + \sum_{\ell=2}^{N_M} \dot{\mathbf{B}}_\ell(t) \mathbf{C}_\ell^M + \sum_{m=1}^{N_T} \dot{\mathbf{B}}_m^T(t) \mathbf{C}_m^T, \quad \mathbf{E}_0 \in \mathbf{H}_0(\mathbf{curl}), \\ \Phi \in H^1(\Omega).$$

\mathbf{E} -based variational problem:

$$\int_{\Omega} \partial_t(\epsilon \mathbf{E})(t) \cdot \mathbf{E}'_0 - \mathbf{H}(t) \cdot \mathbf{curl} \mathbf{E}'_0 \, d\mathbf{x} = 0 \quad \forall \mathbf{E}'_0 \in \mathbf{H}_0(\mathbf{curl}, \Omega),$$

$$\int_{\Omega} \partial_t(\epsilon \mathbf{E})(t) \cdot \mathbf{grad} \Phi', \gamma_p \Phi \in \mathcal{S}(\Gamma) \, d\mathbf{x} = 0 \quad \forall \Phi',$$

$$\int_{\Omega} \partial_t(\epsilon \mathbf{E})(t) \cdot \mathbf{C}_k^E \, d\mathbf{x} - \mathbf{J}_k(t) = 0 \quad \forall k = 2, \dots, N_E,$$

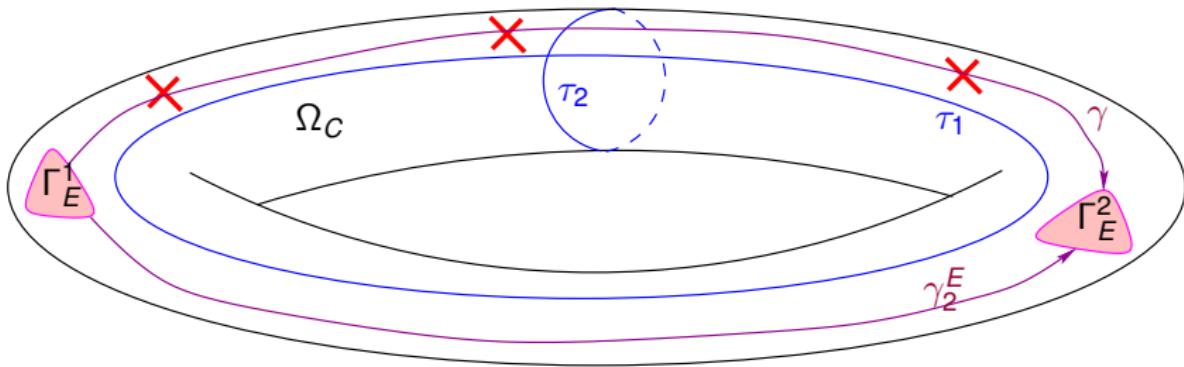
$$\int_{\Omega} \partial_t(\epsilon \mathbf{E})(t) \cdot \mathbf{C}_\ell^M - \mathbf{H}(t) \cdot \mathbf{curl} \mathbf{C}_\ell^E \, d\mathbf{x} - \mathbf{F}_\ell(t) = 0 \quad \forall \ell = 2, \dots, N_M,$$

$$\int_{\Omega} \partial_t(\epsilon \mathbf{E})(t) \cdot \mathbf{C}_m^T - \mathbf{H}(t) \cdot \mathbf{curl} \mathbf{C}_m^T \, d\mathbf{x} - \mathbf{J}_m^T(t) = 0 \quad \forall m = 1, \dots, N_T,$$

$$\int_{\Omega} (\partial_t(\mu \mathbf{H})(t) + \mathbf{curl} \mathbf{E}(t)) \cdot \mathbf{H}' \, d\mathbf{x} = 0 \quad \forall \mathbf{H}' \in \mathbf{L}^2(\Omega).$$

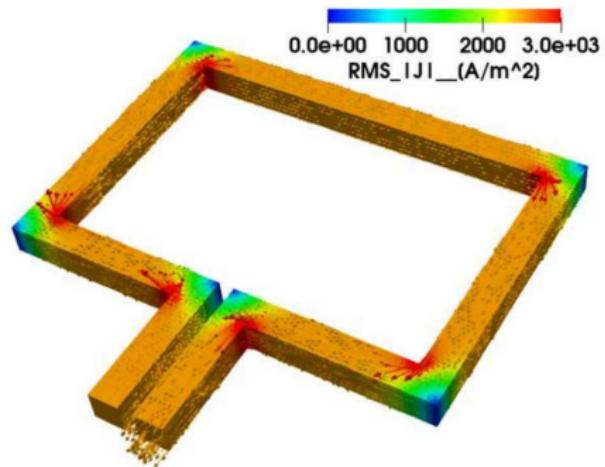
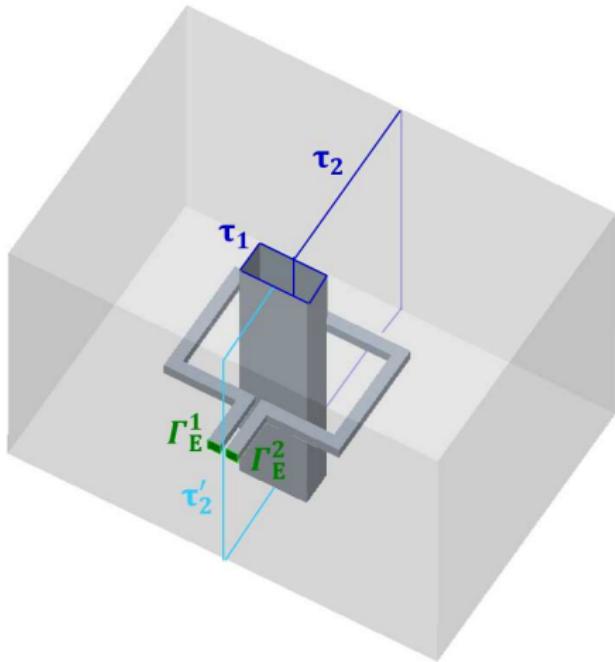
Topological Obstructions

$\beta_1(\Omega_C) > 0 \rightarrow$ unique voltages requires **cuts** in Ω_C !
bounded by topological cycles $\hat{\tau}_m$



The Importance of Choosing Cuts

Computations in frequency domain, bounded Ω , voltage excitation.

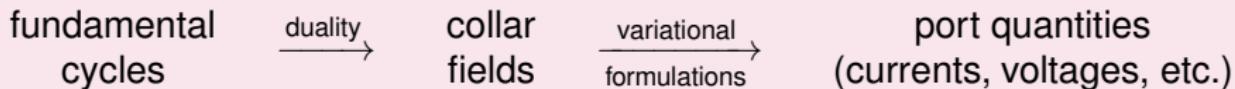


Wrap-Up: Circuit-Field Coupling & Topology

Mathematical tools: Relative (co-)homology,
Poincaré-Lefschetz duality



R. HIPTMAIR AND J. OSTROWSKI, *Electromagnetic port boundary conditions: Topological and variational perspective*, Tech. Rep. 2020-27, Seminar for Applied Mathematics, ETH Zürich, Switzerland, 2020.



$\beta_1(\Omega_c) > 0$ dangerous!